

**Physics 842**  
**Particle Physics, Winter Term 2009**  
**Midterm I**

Thursday Feb 12th 2009, 9:30 am.

Professor Meadows

[To be handed in by 5 pm Monday Feb 16, 2009]

Name .....

**Solution to Question 1**

- a) Write down the parity transformed Dirac spinor corresponding to

$$u(p) = \frac{Nc}{E + mc^2} \begin{pmatrix} E/c + mc \\ 0 \\ p_z \\ p_x + ip_y \end{pmatrix}$$

(Hint: What operator is used for  $P$  on Dirac spinors?)

**Solution:**

We follow the rule that a parity-transformed wave-function  $\psi'$  from the Dirac equation is related to its un-transformed state  $\psi$  by

$$\psi' = \gamma^0 \psi$$

From this we obtain, the present case:

$$u'(p) = \gamma^0 u(p) = \frac{Nc}{E + mc^2} \begin{pmatrix} E/c + mc \\ 0 \\ -p_z \\ -(p_x + ip_y) \end{pmatrix}$$

NOTE that the factor

$$e^{-\frac{i}{\hbar} p_\mu x^\mu} = e^{-\frac{i}{\hbar} (Et - \vec{p} \cdot \vec{r})}$$

does NOT change when  $(x, y, z) \rightarrow (-x, -y, -z)$ .

[15 points]

b) Show that the spinor for an electron at rest is an eigenstate of parity.

**Solution:**

See below.

[15 points]

c) Show that the spinor for a positron at rest is also an eigenstate of parity.

**Solution:**

A positron at rest has spinor has

$$\text{either } v^+(p) = N \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} \text{ (positive helicity) or}$$
$$v^-(p) = N \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} \text{ (negative helicity).}$$

The parity operation on each of these yields:

$$\text{either } v^{+'} = \gamma^0 v^+(p) = N \begin{pmatrix} 0 \\ 0 \\ -1 \\ 0 \end{pmatrix} = -v^+(p) \text{ (positive helicity) or}$$
$$v^{+'} = \gamma^0 v^-(p) = N \begin{pmatrix} 0 \\ 0 \\ 0 \\ -1 \end{pmatrix} = -v^-(p) \text{ (negative helicity).}$$

[15 points]

d) How are the parity of the electron and positron related?

**Solution:**

An electron at rest has spinor

$$\text{either } u^+(p) = N \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \text{ (positive helicity) or}$$

$$u^-(p) = N \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \text{ (negative helicity).}$$

The parity operation on each of these yields:

$$\text{either } u^{+'} = \gamma^0 u^+(p) = N \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = +u^+(p) \text{ (positive helicity) or}$$

$$u^{+'} = \gamma^0 u^-(p) = N \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} = +u^-(p) \text{ (negative helicity).}$$

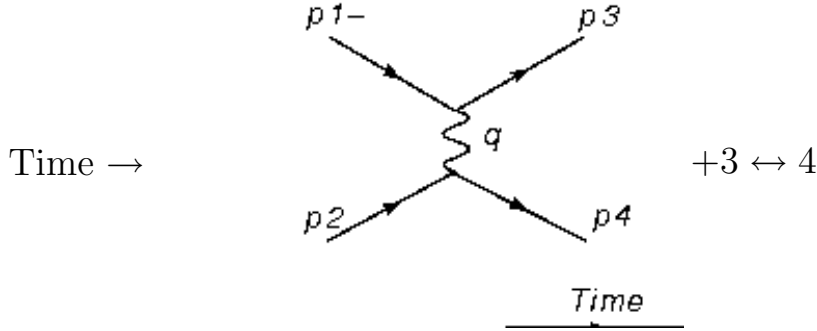
From part (c) it is clear that, with this convention, the positron has parity  $P_{e^+} = -1$  and the electron has  $P_{e^-} = +1$ . In any convention (we can choose  $P = -\gamma^0$  rather than  $P = \gamma^0$ ) the  $e^-$  and  $e^+$  have opposite parity.

[15 points]

**Solution to Question 2**

Sketch the two lowest order diagrams contributing to  $e^-e^- \rightarrow e^-e^-$  (Moller) scattering.

**Solution:**



[15 points]

- a) Use the Feynman rules to write down the matrix element  $\mathcal{M}_1$  for the first diagram you sketched.

**Solution:**

$$\begin{aligned}
 -i\mathcal{M}_1 &= (2\pi)^4 \int [\bar{u}(3)ig_e\gamma^\mu u(1)] \frac{-ig_{\mu\nu}}{q^2} [\bar{u}(4)ig_e\gamma^\nu u(2)] \\
 &\quad \times \delta^4(p_1 - p_3 - q)\delta^4(p_2 + q - p_4)d^4q \\
 \mathcal{M}_1 &= -\frac{g_e^2}{q^2} [\bar{u}(3)\gamma^\mu u(1)] [\bar{u}(4)\gamma_\mu u(2)] \\
 \text{where } q &= (p_1 - p_3)
 \end{aligned}$$

[15 points]

- b) Use the Feynman rules to write down the matrix element  $\mathcal{M}_2$  for the second diagram.

**Solution:**

$$\begin{aligned}
 -i\mathcal{M}_2 &= (2\pi)^4 \int [\bar{u}(4)ig_e\gamma^\mu u(1)] \frac{-ig_{\mu\nu}}{q^2} [\bar{u}(3)ig_e\gamma^\nu u(2)] \\
 &\quad \times \delta^4(p_1 - p_4 - q)\delta^4(p_2 + q - p_3)d^4q
 \end{aligned}$$

$$\mathcal{M}_2 = -\frac{g_e^2}{q^2} [\bar{u}(4)\gamma^\mu u(1)] [\bar{u}(3)\gamma_\mu u(2)]$$

where  $q = (p_1 - p_4)$

[15 points]

- c) Write  $\mathcal{M}$ , as the correct combination for these two matrix elements.

**Solution:**

Since the electrons in the final state are identical, we need to antisymmetrize wrt their interchange. So we get

$$\mathcal{M} = \mathcal{M}_1 - \mathcal{M}_2$$

[15 points]

- d) Expand  $|\mathcal{M}|^2$ , as a sum of terms in Dirac spinors.

**Solution:**

$$\begin{aligned} |\mathcal{M}|^2 &= |\mathcal{M}_1|^2 + |\mathcal{M}_2|^2 - \mathcal{M}_1\mathcal{M}_2^* - \mathcal{M}_2\mathcal{M}_1^* \\ &= \frac{g_e^4}{(p_1 - p_3)^4} |[\bar{u}(3)\gamma^\mu u(1)] [\bar{u}(4)\gamma_\mu u(2)]|^2 \\ &\quad + \frac{g_e^4}{(p_1 - p_4)^4} |[\bar{u}(4)\gamma^\mu u(1)] [\bar{u}(3)\gamma_\mu u(2)]|^2 \\ &\quad - \frac{g_e^4}{(p_1 - p_4)^2 (p_1 - p_3)^2} \\ &\quad \times \{ [\bar{u}(3)\gamma^\mu u(1)] [\bar{u}(4)\gamma_\mu u(2)] [\bar{u}(4)\gamma^\mu u(1)]^* [\bar{u}(3)\gamma_\mu u(2)]^* \\ &\quad + [\bar{u}(4)\gamma^\mu u(1)] [\bar{u}(3)\gamma_\mu u(2)] [\bar{u}(3)\gamma^\mu u(1)]^* [\bar{u}(4)\gamma_\mu u(2)]^* \} \end{aligned}$$

[15 points]

### Solution to Question 3

For  $e^-e^- \rightarrow e^-e^-$  scattering introduced above:

- a) Use “Casimir’s trick” to evaluate the matrix element  $\langle |\mathcal{M}_1|^2 \rangle$  summed over spins in the final state and averaged over the initial spins.

**Solution:**

“Casimir’s trick” sums over spins in the final state and also over initial state spins. In short, it converts a term like

$$\sum_{s_a, s_b} [\bar{u}^{s_a}(p_a)\Gamma_1 u^{s_b}(p_b)] [\bar{u}^{s_a}(p_a)\Gamma_2 u^{s_b}(p_b)]^*$$

where  $\Gamma_{1,2}$  are 4 x 4 matrices to

$$Tr [\Gamma_1(\not{p}_b + m_b c)\bar{\Gamma}_2(\not{p}_a + m_a c)]$$

where  $\bar{\Gamma}_2 = \gamma^0 \Gamma_2^\dagger \gamma^0$ .

Applying this to  $\mathcal{M}_1$  where  $\Gamma_2 = \gamma^\nu$  and  $\bar{\Gamma}_2 = \gamma^0 \gamma^{\nu\dagger} \gamma^0 = \gamma^\nu$ :

$$\begin{aligned} \mathcal{M}_1 &= -\frac{g_e^2}{(p_1 - p_3)^2} [\bar{u}(3)\gamma^\mu u(1)] [\bar{u}(4)\gamma_\mu u(2)] \\ \rightarrow \langle |\mathcal{M}_1|^2 \rangle &= \frac{g_e^4}{4(p_1 - p_3)^4} Tr [\gamma^\mu(\not{p}_1 + mc)\gamma^\nu(\not{p}_3 + mc)] \\ &\quad \times Tr [\gamma_\mu(\not{p}_2 + mc)\gamma_\nu(\not{p}_4 + mc)] \end{aligned}$$

[For the other diagram, simply exchange  $p_3 \leftrightarrow p_4$ .] The factor 4 is equal to the number of initial spin states,  $(2s_a + 1)(2s_b + 1)$ , and is required to convert the sum to an average. [15 points]

- b) Evaluate this as a function of the 4-momenta  $p_{1-4}$  of the ingoing and outgoing electrons.

**Solution:**

Expanding the first term:

$$Tr (\gamma^\mu \not{p}_1 \gamma^\nu \not{p}_3) + mc [Tr (\gamma^\mu \not{p}_1 \gamma^\nu) + Tr (\gamma^\mu \gamma^\nu \not{p}_3)] m^2 c^2 Tr (\gamma^\nu \gamma^\nu)$$

Using

$$Tr \{\text{prod. of odd number of } \gamma\text{'s}\} = 0,$$

the terms in  $mc$  are zero. Using

$$\text{Tr}\{\gamma^\alpha\gamma^\beta\gamma^\gamma\gamma^\delta\} = 4(g^{\alpha\beta}g^{\gamma\delta} - g^{\alpha\gamma}g^{\beta\delta} + g^{\alpha\delta}g^{\beta\gamma})$$

the first term becomes

$$\begin{aligned} \text{Tr}(\gamma^\mu\not{p}_1\gamma^\nu\not{p}_3) &= p_{1\alpha}p_{3\beta}\text{Tr}(\gamma^\mu\gamma^\alpha\gamma^\nu\gamma^\beta) \\ &= p_{1\alpha}p_{3\beta}4(g^{\mu\alpha}g^{\nu\beta} - g^{\mu\nu}g^{\alpha\beta} + g^{\mu\beta}g^{\alpha\nu}) \\ &= 4(p_1^\mu p_3^\nu - g^{\mu\nu}(p_1 \cdot p_3) + p_3^\mu p_1^\nu) \end{aligned}$$

Finally, for the  $m^2c^2$  term:

$$\text{Tr}(\gamma^\alpha\gamma^\beta) = 4g^{\mu\nu}$$

For the first factor in square parentheses, therefore, we obtain

$$\text{Tr}[\gamma^\mu(\not{p}_1 + mc)\gamma^\nu(\not{p}_3 + mc)] = 4[p_1^\mu p_3^\nu + p_3^\mu p_1^\nu + g^{\mu\nu}(m^2c^2 - p_1 \cdot p_3)]$$

The second factor in square parentheses is similar ( $p_1 \rightarrow p_2$  and  $p_3 \rightarrow p_4$ ) leading to

$$\begin{aligned} \langle |\mathcal{M}_1|^2 \rangle &= \frac{4g_e^4}{(p_1 - p_3)^4} [p_1^\mu p_3^\nu + p_3^\mu p_1^\nu + g^{\mu\nu}(m^2c^2 - p_1 \cdot p_3)] \\ &\quad \times [p_{2\mu}p_{4\nu} + p_{4\mu}p_{2\nu} + g_{\mu\nu}(m^2c^2 - p_2 \cdot p_4)] \\ &= \frac{4g_e^4}{(p_1 - p_3)^4} [2(p_1 \cdot p_2)(p_3 \cdot p_4) + 2(p_1 \cdot p_4)(p_2 \cdot p_3) \\ &\quad - 2(m^2c^2)(p_2 \cdot p_4 + p_1 \cdot p_3) + 4(m^2c^2)^2] \end{aligned}$$

[For the other diagram, simply exchange  $p_3 \leftrightarrow p_4$ .] [15 points]

- c) Evaluate  $\langle |\mathcal{M}_1|^2 \rangle$  for scattering in the CMS where the incoming momenta have magnitude  $|\vec{p}| \gg mc$ .

**Solution:**

In this frame

$$\begin{aligned} p_1 &= (E/c, 0, 0, p) ; p_2 = (E/c, 0, 0, -p) ; \\ p_3 &= (E/c, p \sin \theta, 0, p \cos \theta) ; p_4 = (E/c, -p \sin \theta, 0, -p \cos \theta) \end{aligned}$$

If, further  $|\vec{p}| \gg mc$ , then  $E/c \approx p$  so that

$$\begin{aligned}
 (p_1 - p_3)^2 &= -2p^2(1 - \cos \theta) = -4p^2 \sin^2 \frac{\theta}{2} \\
 (p_1 - p_4)^2 &= -2p^2(1 + \cos \theta) = -4p^2 \cos^2 \frac{\theta}{2} \quad \text{and} \\
 p_1 \cdot p_2 &= 2p^2 ; p_1 \cdot p_3 = p^2(1 - \cos \theta) ; p_1 \cdot p_4 = p^2(1 + \cos \theta) ; \\
 p_3 \cdot p_4 &= 2p^2 ; p_2 \cdot p_3 = p^2(1 + \cos \theta) ; p_2 \cdot p_4 = p^2(1 - \cos \theta)
 \end{aligned}$$

Insert this into the expression for  $\langle |\mathcal{M}_1|^2 \rangle$ , neglecting the terms in  $m^2 c^2$  in comparison with  $p^2$ :

$$\begin{aligned}
 \langle |\mathcal{M}_1|^2 \rangle &= -\frac{g_e^4}{4 \sin^4(\theta/2)} [8 + 2(1 + \cos \theta)(1 + \cos \theta)] \\
 &= -\frac{g_e^4}{2 \sin^4(\theta/2)} [5 + 2 \cos \theta + \cos^2 \theta]
 \end{aligned}$$

For the other diagram, simply exchange  $p_3 \leftrightarrow p_4$ . This will give

$$\begin{aligned}
 \langle |\mathcal{M}_2|^2 \rangle &= -\frac{g_e^4}{4 \cos^4(\theta/2)} [8 + 2(1 - \cos \theta)(1 - \cos \theta)] \\
 &= -\frac{g_e^4}{2 \cos^4(\theta/2)} [5 - 2 \cos \theta + \cos^2 \theta]
 \end{aligned}$$

[15 points]