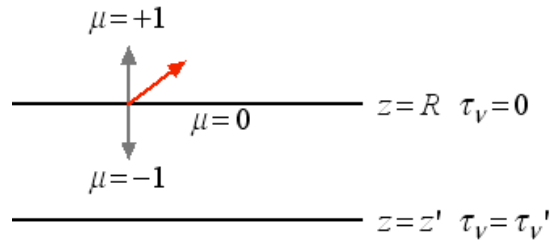


10 - GENERAL SOLUTION OF THE EQUATION OF TRANSFER

$$\frac{dI_v}{dz \sec \theta} = \cos \theta \frac{dI_v}{dz} = j_v \rho - k_v \rho I_v$$



It is conventional to set $\tau = 0$ at $z = R$ (the "surface" of the star), and have τ increase inward as z decreases.

$$\text{i.e., } d\tau_v = -k_v \rho dz, \quad \tau_v = \int_R^{z'} -k_v \rho dz$$

We will also use $\mu = \cos \theta$ with $\mu = 1$ pointing outward from the stellar surface.

$$\mu \frac{dI_v}{dz} = j_v \rho - k_v \rho I_v \quad \text{now } \div \text{ by } k_v \rho,$$

$$\mu \frac{dI_v}{d\tau_v} = I_v - \frac{j_v}{k_v} = I_v - S_v$$

$$\text{where } S_v \equiv \frac{j_v}{k_v} \quad \text{the "Source Function"}$$

$$[\text{In TE, } S_v = B_v \quad \text{and} \quad j_v = k_v B_v]$$

EQUATION OF (RADIATIVE) TRANSFER FOR STARS

$$\mu \frac{dI_v}{d\tau_v} = I_v - S_v \quad \text{"The Equation"}$$

Let's divide by μ , move I_v over to the other side,

$$\frac{dI_v}{d\tau_v} - \frac{I_v}{\mu} = -\frac{S_v}{\mu}$$

For a general equation of the form $\frac{dy}{dx} + f(x)y = g(x)$ we can solve by multiplying

through by $e^{\int f(x)dx}$ to get

$$\left(\frac{dI_v}{d\tau_v} - \frac{I_v}{\mu} \right) e^{-\tau_v/\mu} = -\frac{S_v}{\mu} e^{-\tau_v/\mu}$$

To solve this, we need some boundary conditions

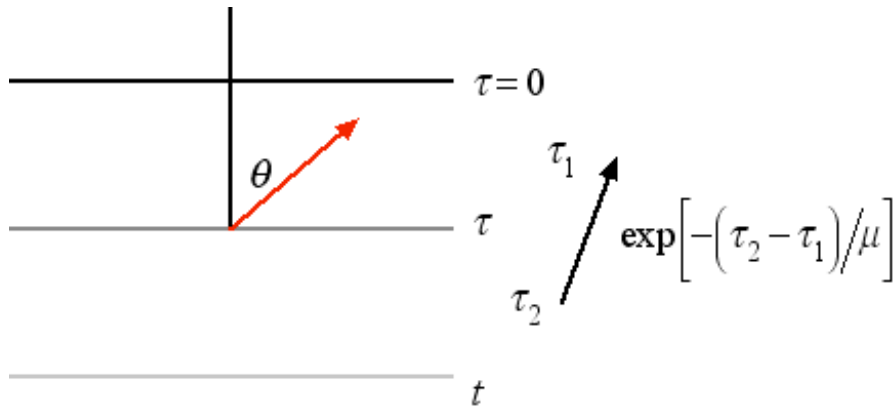
1. For outward flowing radiation:

$$\lim_{\tau \rightarrow \infty} I_v(\tau_v, \mu) e^{-\tau_v/\mu} \rightarrow 0$$

2. For inward flowing radiation:

$$I_v(0, \mu) = 0 \quad \text{for } -1 \leq \mu \leq 0$$

Let us drop the v subscript for now (saves typing!)



Let t be the variable, and τ be the limits of the problem.

$$\frac{d}{dt} \left(I e^{-t/\mu} \right) = -\frac{S e^{-t/\mu}}{\mu}$$

$$I e^{-t/\mu} \Big|_{t=\tau_1}^{t=\tau_2} = -\int_{\tau_1}^{\tau_2} S(t) e^{-t/\mu} \frac{dt}{\mu} \quad \text{Then } \times e^{-\tau_1/\mu} \text{ to get :}$$

$$I(\tau_1, \mu) = I(\tau_2, \mu) e^{-(\tau_2 - \tau_1)/\mu} + \int_{\tau_1}^{\tau_2} S(t) e^{-(t - \tau_1)/\mu} \frac{dt}{\mu}$$

Example \Rightarrow Emergent intensity for a semi-infinite atmosphere: Take $\tau_1 = 0, \tau_2 \rightarrow \infty$

$$I(0, \mu) = \int_0^{\infty} S(t) e^{-t/\mu} \frac{dt}{\mu}$$

The emergent intensity is a weighted mean over the source function, and corresponds to the integral of the fraction of the radiant energy at each layer, attenuated by extinction.

[NOTE for the mathematically-minded: I is the Laplace Transform of S]

Now, let us look at some simple, but illustrative examples.

If (for whatever reason) S is linear with optical depth, $S = a + bt$, then

$$I(0, \mu) = a + b\mu$$

For a finite slab of thickness τ , with no incident radiation, the emergent I is, for $S = \text{constant}$,

$$I(0, 1) = S(1 - e^{-\tau})$$

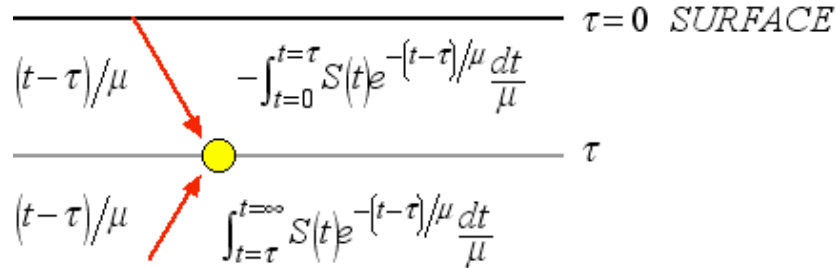
For $\tau \gg 1$ (optically thick case),

$$I = S$$

For $\tau \ll 1$ (optically thin case)

$$I \approx S\tau \quad (\text{where } e^{-\tau} \approx 1 - \tau \text{ for } \tau \ll 1)$$

Arbitrary Point ($\tau_1 = \tau$) in a Semi-Infinite Atmosphere with No Incident Radiation.



FOR $\mu \geq 0$ (OUTGOING RADIATION) ($\tau_2 \rightarrow \infty$)

$$I(\tau, \mu) = \lim_{\tau_2 \rightarrow \infty} I(\tau_2, \mu) e^{-(\tau_2 - \tau)/\mu} + \int_{\tau}^{\infty} S(t) e^{-(t - \tau)/\mu} \frac{dt}{\mu}$$

$$I(\tau, \mu) = \int_{\tau}^{\infty} S(t) e^{-(t - \tau)/\mu} \frac{dt}{\mu} \quad (0 \leq \mu \leq 1)$$

FOR $\mu \leq 0$ (INCOMING RADIATION) ($\tau_2 \rightarrow 0$)

$$I(\tau, \mu) = I(0, \mu) e^{\tau/\mu} + \int_{\tau}^0 S(t) e^{-(t - \tau)/\mu} \frac{dt}{\mu}$$

$$I(\tau, \mu) = -\int_0^{\tau} S(t) e^{-(\tau - t)/\mu} \frac{dt}{\mu} \quad (-1 \leq \mu \leq 0)$$

These are the "I Integrals"

J INTEGRAL

$$\begin{aligned}
 J_v &= \frac{1}{2} \int_{-1}^1 I_v d\mu \\
 &= \frac{1}{2} \left[\int_0^1 \int_{\tau_v}^{\infty} S_v e^{-(t_v - \tau_v)/\mu} dt_v \frac{d\mu}{\mu} + \int_{-1}^0 (-) \int_0^{\tau_v} S(t) e^{(\tau_v - t_v)/\mu} dt_v \frac{d\mu}{\mu} \right]
 \end{aligned}$$

After some algebra that I won't bore you with,

$$J_v = \frac{1}{2} \left\{ \int_{\tau_v}^{\infty} S_v E_1(t_v - \tau_v) dt_v + \int_0^{\tau_v} S_v E_1(\tau_v - t_v) dt_v \right\}$$

Where we have defined the Exponential integral of degree n as

$$E_n(x) = \int_1^{\infty} \frac{e^{-wx}}{w^n} dw$$

H INTEGRAL

$$H_v = \frac{1}{2} \left\{ \int_{\tau_v}^{\infty} S_v E_2(t_v - \tau_v) dt_v - \int_0^{\tau_v} S_v E_2(\tau_v - t_v) dt_v \right\}$$

K INTEGRAL

$$K_v = \frac{1}{2} \left\{ \int_{\tau_v}^{\infty} S_v E_3(t_v - \tau_v) dt_v + \int_0^{\tau_v} S_v E_3(\tau_v - t_v) dt_v \right\}$$

Useful relations for $E_n(x)$

$$E_n(x) = \frac{1}{n-1} \left[e^{-x} - x E_{n-1}(x) \right] \quad n \geq 2$$

$$E_1(x) = \frac{e^{-x}}{x} \left[1 - \frac{1}{x} + \frac{2!}{x^2} - \frac{3!}{x^3} \dots \right] \text{ for } x \gg 1$$

$$E_1(x) = \gamma - \ln x + \sum_{k=1}^{\infty} (-1)^{k-1} \frac{x^k}{kk!} \text{ for } x > 0$$

($\gamma = 0.5572156\dots$ Euler's constant)